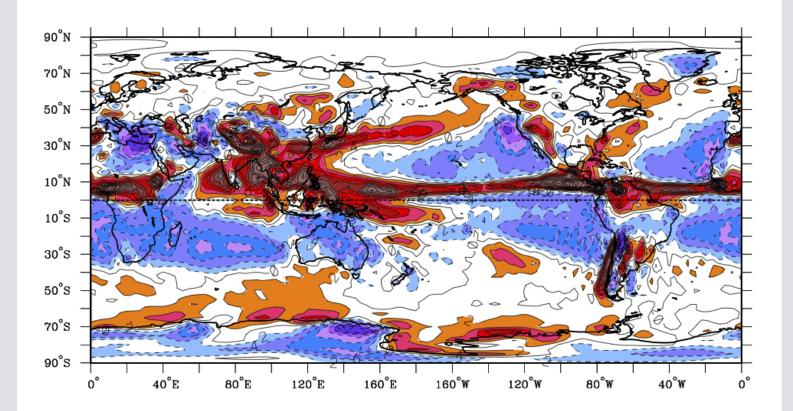
June 40 year mean vertical velocity



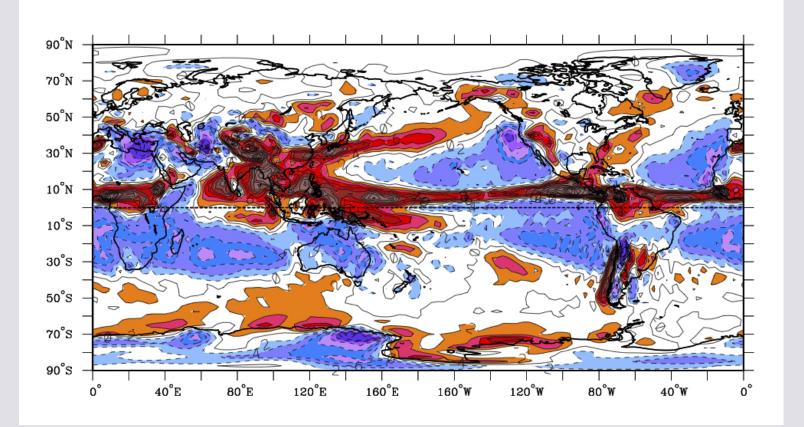
PDEs and Asymptotics for the Tropical Atmosphere

Joseph A. Biello

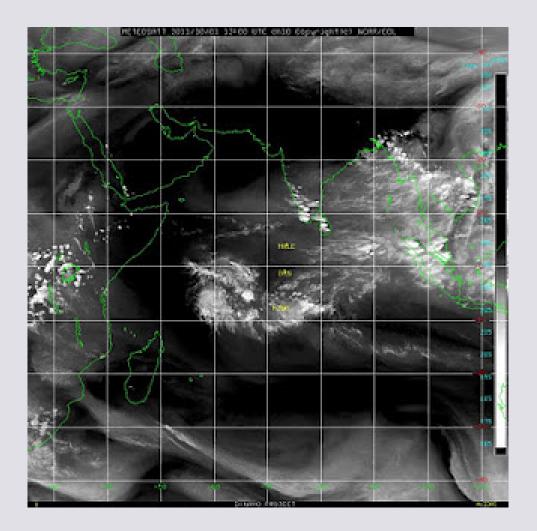
- joint work with Andy Majda (NYU), Francesco Rubini (U. Florence), Paolo Ruggieri (U. Bologna), Xiaoyun Niu (UC Davis)
- Multiple scales wave/mean flow theory of the Madden-Julian oscillation
- Interpreting the results of the ECMWF climate model using the multi-scale theory.



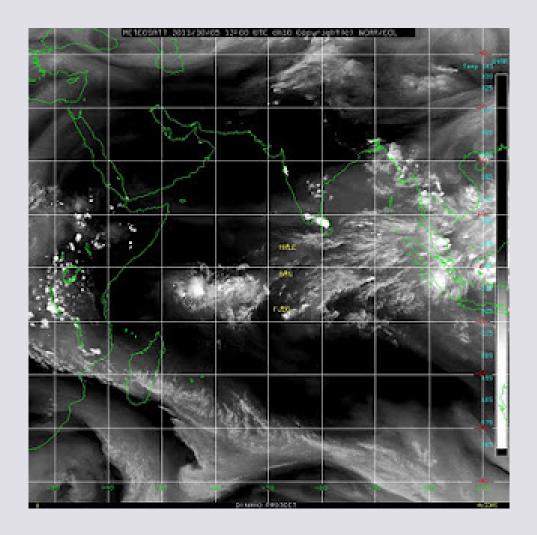
June 40 year mean vertical velocity



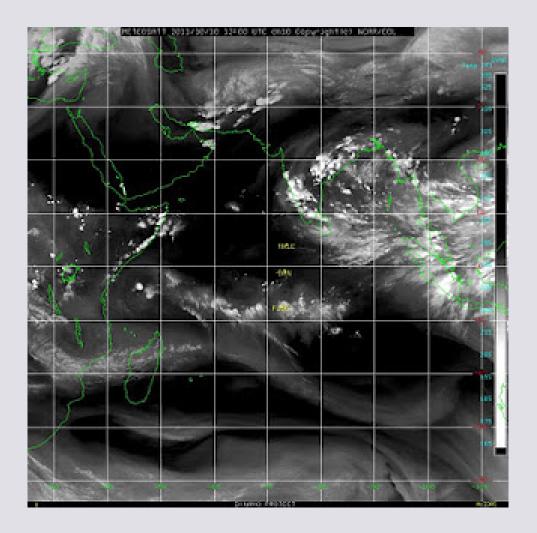
MJO in the Infrared: DYNAMO , Oct 1, 2011



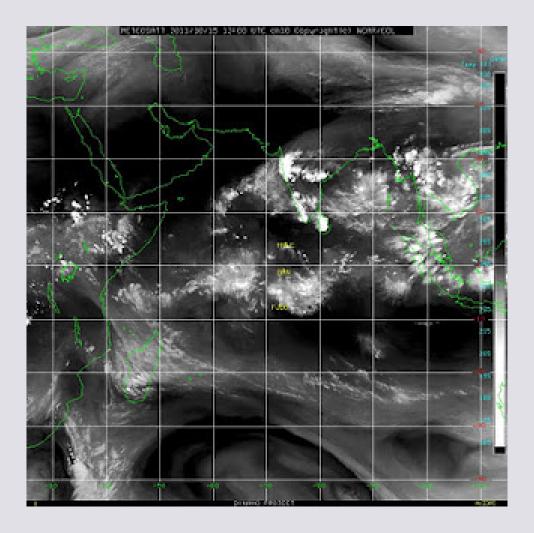
MJO in the Infrared: DYNAMO , Oct 5, 2011



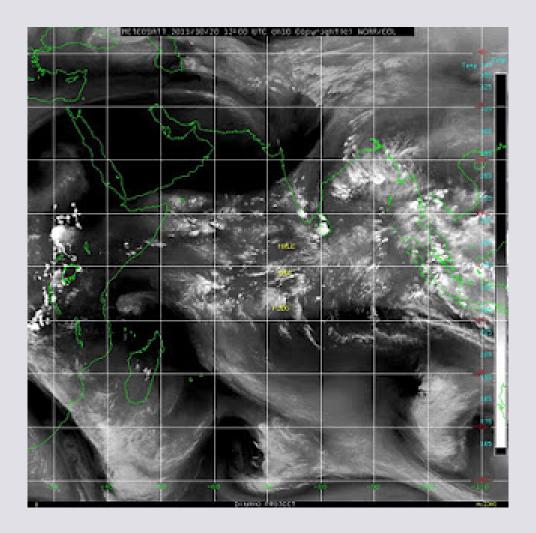
MJO in the Infrared: DYNAMO , Oct 10, 2011



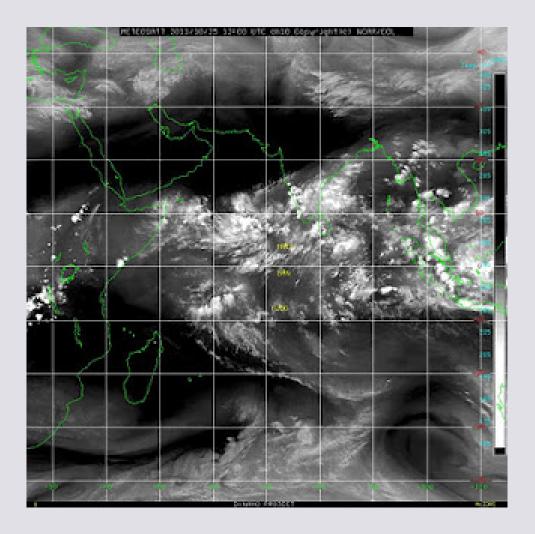
MJO in the Infrared: DYNAMO , Oct 15, 2011



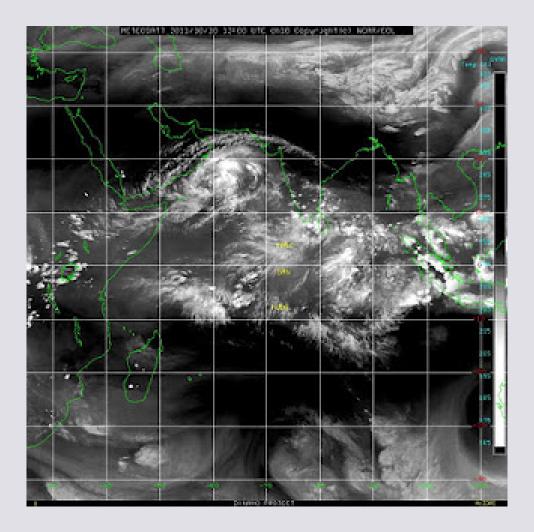
MJO in the Infrared: DYNAMO , Oct 20, 2011



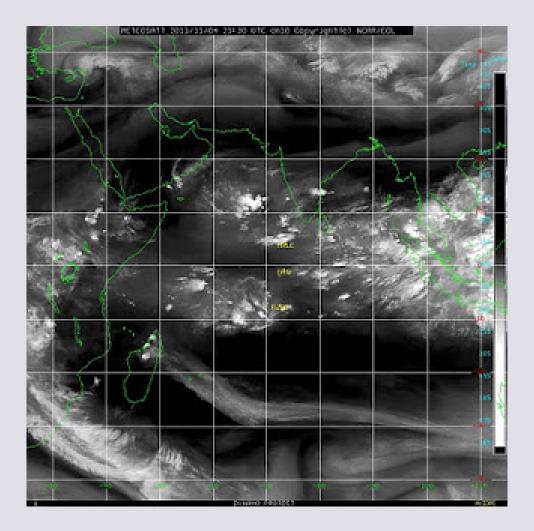
MJO in the Infrared: DYNAMO , Oct 25, 2011



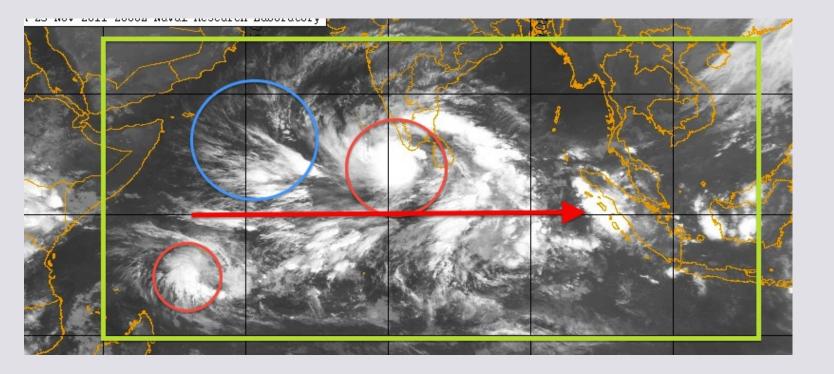
MJO in the Infrared: DYNAMO , Oct 30, 2011



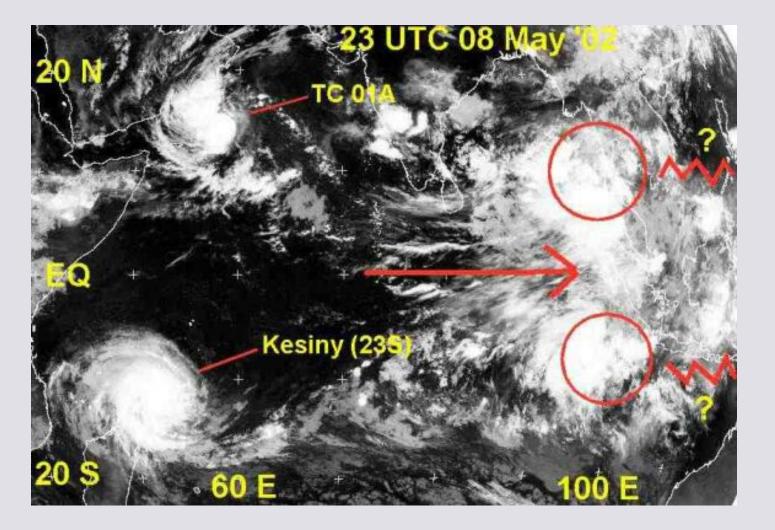
MJO in the Infrared: DYNAMO , Nov 4, 2011



MJO in visible light



MJO in the Infrared

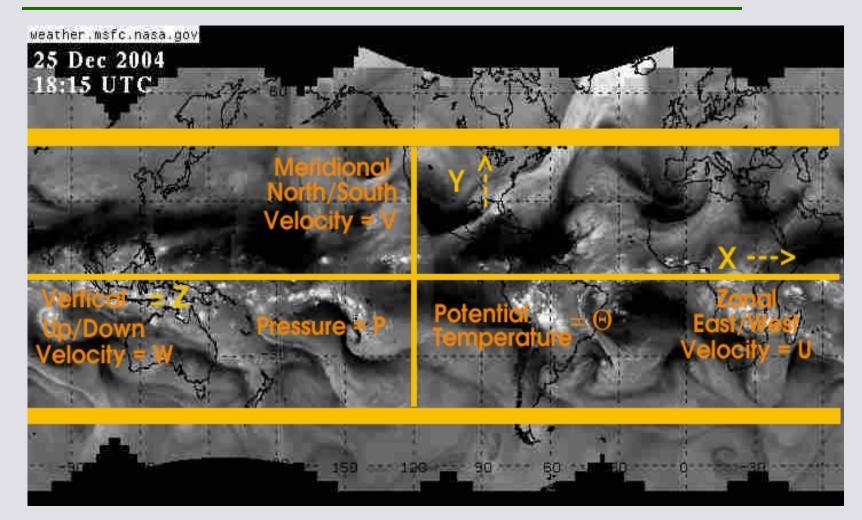


MJO in the Infrared

Overview of the Tropical Atmosphere

- Troposphere is extremely thin (16 : 40000 km)
- Length scales range from rain drop size to planetary scale.
- Time scales range from fractions of an hour to days to a few weeks (intraseasonal) to climate timescales.
- Coriolis parameter vanishes at the equator causing equatorial waveguide and unique nonlinear dynamics.
- Nonlinearities are strong on small scales, but weak on the largest scales.
- Forcing is dominated by latent heat release from water condensation when clouds are formed.

Coordinate System



The Equatorial Primitive Equations

Incompressible, Hydrostatic, Coriolis force on β -plane

 $u_t + \vec{u} \cdot \nabla u - \beta y v = -p_x + S_u$ $v_t + \vec{u} \cdot \nabla v + \beta y u = -p_y + S_v$ $\theta_t + \vec{u} \cdot \nabla \theta + N^2 w = S_\theta$ $p_z = \theta$ $u_x + v_y + w_z = 0$

- $\vec{u} = (u, v, w) = (\text{East, North, Up})$
- $\theta = \text{potential temperature perturbation}$
- p = pressure perturbation

•
$$N = \sqrt{\frac{g}{\theta_0} \frac{d\theta_0}{dz}} =$$
 buoyancy frequency

- βy is vertical component of Coriolis force near equator
- Rigid lid $\Rightarrow w = 0$ at $z = 0, 16 \,\mathrm{km}$
- Heat and momentum sources and sinks $= S_{\theta}, S_u, S_v$

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Incompressible, Hydrostatic, Coriolis force on β -plane

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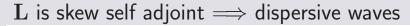
•
$$N = \sqrt{\frac{g}{\theta_0} \frac{d\theta_0}{dz}} =$$
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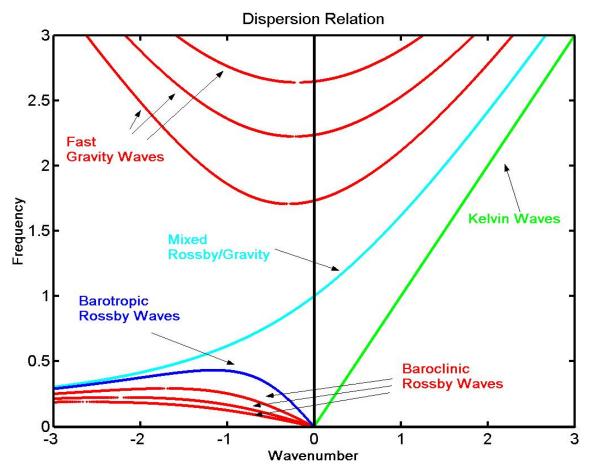
- βy is vertical component of Coriolis force near equator
- Rigid lid $\Rightarrow w = 0$ at $z = 0, 16 \,\mathrm{km}$
- Heat and momentum sources and sinks $= S_{\theta}, S_u, S_v$

Schematically, these are equivalent to

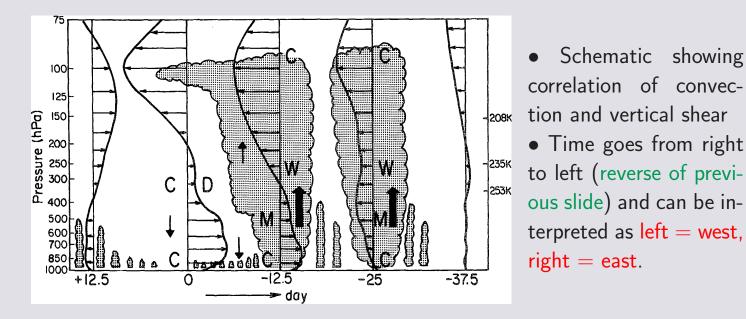
$$\Psi_t + \mathbf{L} \Psi + \mathbf{N} (\Psi, \Psi) = \mathbf{S}(x, t)$$

Linear Theory of Equatorial Waves





MJO: Vertical Shear and Convection



Lin & Johnson J. Atmos. Sci., 53, p 701, fig. 16.

- Congestus clouds weak winds/easterlies
- Westward tilted anvil westerly onset
- Strong westerlies trail convection

Synoptic Scale (Balanced) Dynamics: Planetary Scale Quasi-Linear Dynamics:

$$\begin{aligned} u'_{\tau} - y v' + p'_{x} &= S'_{u} & \overline{U}_{t} - y \overline{V} + \overline{P}_{X} &= F^{U} - d_{0} \overline{U} \\ v'_{\tau} + y u' + p'_{y} &= S'_{v} & y \overline{U} + \overline{P}_{y} &= 0 \\ \theta'_{\tau} + w' &= S'_{\theta} & \overline{\Theta}_{t} + \overline{W} &= F^{\theta} - d_{\theta} \overline{\Theta} + \overline{S}_{\theta} \\ p'_{z} &= \theta' & \overline{P}_{z} &= \overline{\Theta} \\ u'_{x} + v'_{y} + w'_{z} &= 0 & \overline{U}_{X} + \overline{V}_{y} + \overline{W}_{z} &= 0 \\ \overline{S'_{\theta}} &= 0 & \end{aligned}$$

The fluxes from the synoptic scales are given by

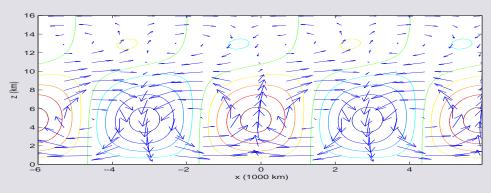
$$\begin{array}{l} F^U &= -\overline{(v'\,u')_y} - \overline{(w'\,u')_z} \\ F^\theta &= -\overline{(v'\,\theta')_y} - \overline{(w'\,\theta')_z} \end{array} \end{array}$$

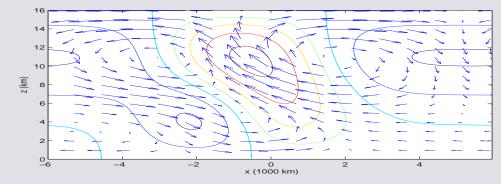
Each forcing effect, i.e. upscale vertical and meridional momentum and temperature transport and planetary scale mean heating can be considered separately and superposed

MJO Model: Convection organized on small scales

- Heating rate traces cloudiness (latent heat release).
- Fluctuations on 1500 km spatial scales
- Clouds/heating localized near equator above Western Pacific.

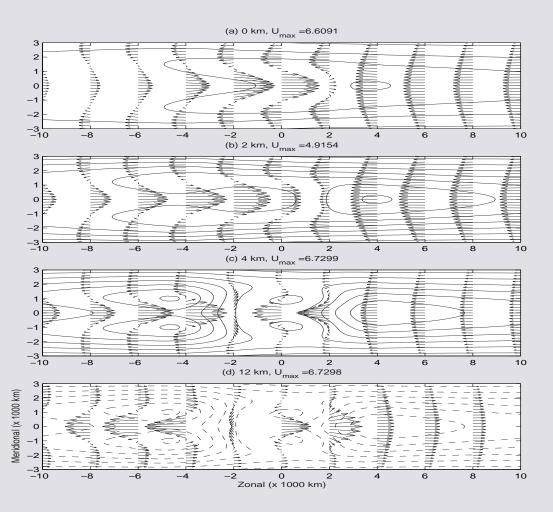
- East: Lower troposphere congestus clouds
- West: High, westward tilted anvil *superclusters*
- Flow vectors and heating contours
- Upscale flux, $\overline{\mathbf{N}(\psi', \psi')} \neq 0$ \Rightarrow Vertical/Longitudinal Tilt





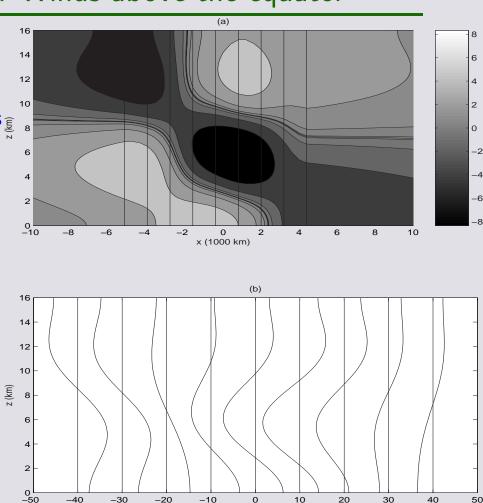
Equatorial MJO model: Flow in the Horizontal Plane

- Congestus heating in the east and westward tilted superclusters in the west of a moving warm pool.
- Planetary mean heating is weaker, but has same structure of synoptic scale fluctuations.
- Pressure and flow at z = 0, 2, 4, 12 km.



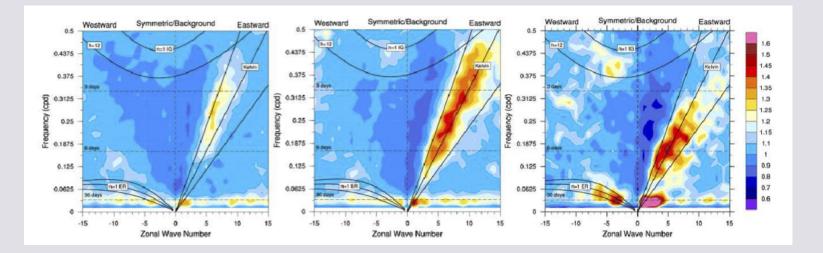
Equatorial MJO model: Winds above the equator

- Lower troposphere congestus heating in the east
- Westward tilted anvil superclusters in the west
- (a) Zonal velocity: westerly = light, easterly = dark versus height and longitude above equator
- (b) Height vs Velocity.

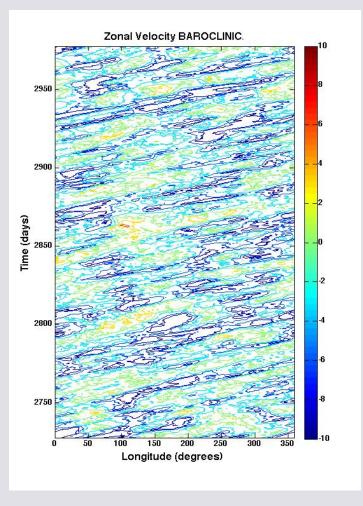


U (m/s)

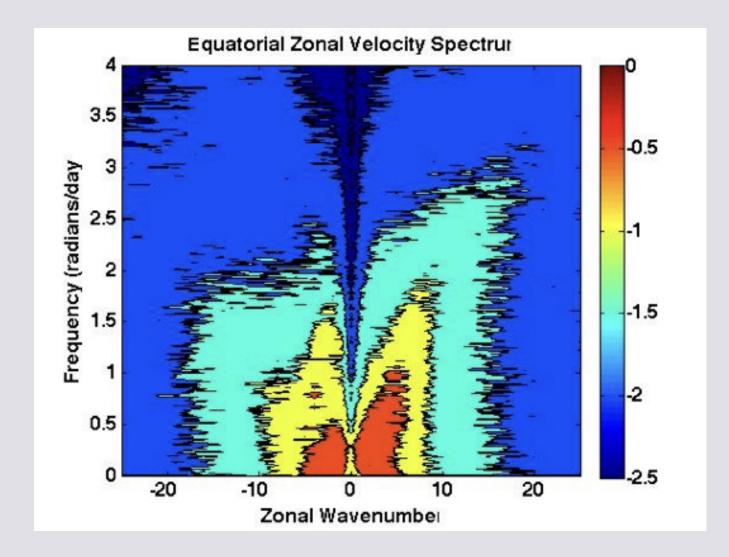
Spectral features of computation and data



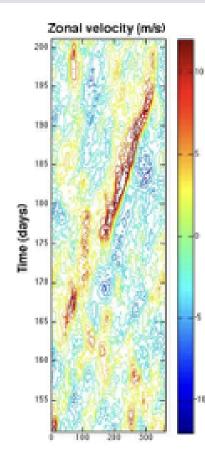
Space time diagram of ECMWF results - Kelvin Waves

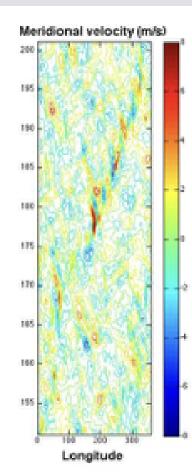


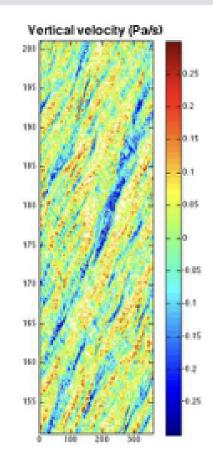
Spectrum of ECMWF results



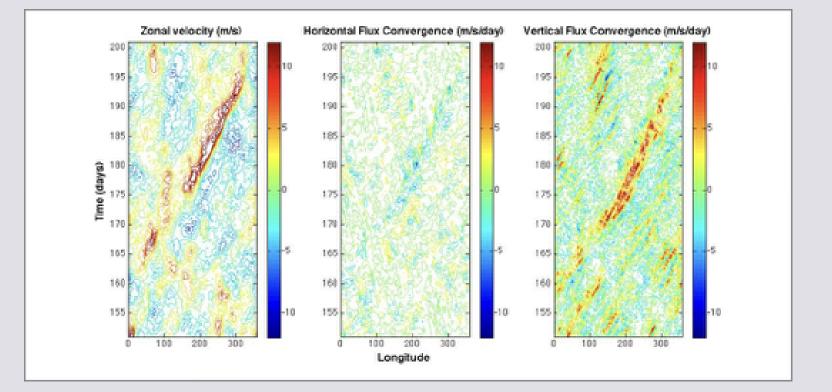
Space time diagram of ECMWF results - Westerly Wind Burst



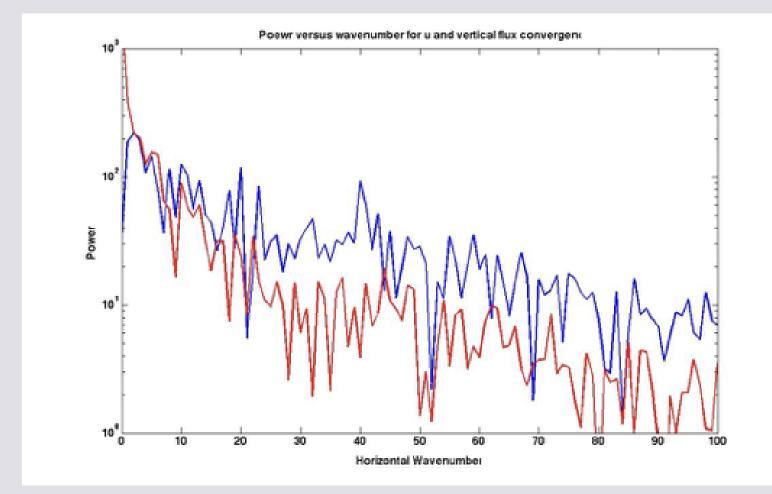




Space time diagram of ECMWF results - Westerly Wind Burst



Fluxes from ECMWF results



In Conclusion

- Numerical simulations are still unable to simulate large scale organized convection.
- I believe that this is due to the inability of the convective models to generate vertical and meridionally tilted structures.

$$\begin{array}{ll} F^U &= -\overline{(v'\,u')_y} - \overline{(w'\,u')_z} \\ F^\theta &= -\overline{(v'\,\theta')_y} - \overline{(w'\,\theta')_z} \end{array}$$

